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The theory of equilibrium reconstruction and a possibility of complete reconstruction in ITER¹

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Nova Photonics

PPPL Experimental Seminar,

March 20, 2007, PPPL, Princeton NJ



Abstract

Potential variances in q- and p- profiles have been calculated for different sets of external and internal measurements envisioned for equilibrium reconstruction in ITER.

It was shown that complementing the external magnetic measurements with either Stark line polarization signals (MSE-LP) or with recently proposed for ITER by Nova Photonics line shift signals (MSE-LS) can significantly improve the reliability of the reconstructed plasma profiles and the magnetic configuration.

Capabilities of calculating variances, incorporated into the numerical code ESC, have completed the theory of reconstruction, which for a long time had a significant gap in ability to evaluate the quality of the presently widely used equilibrium reconstruction technique.

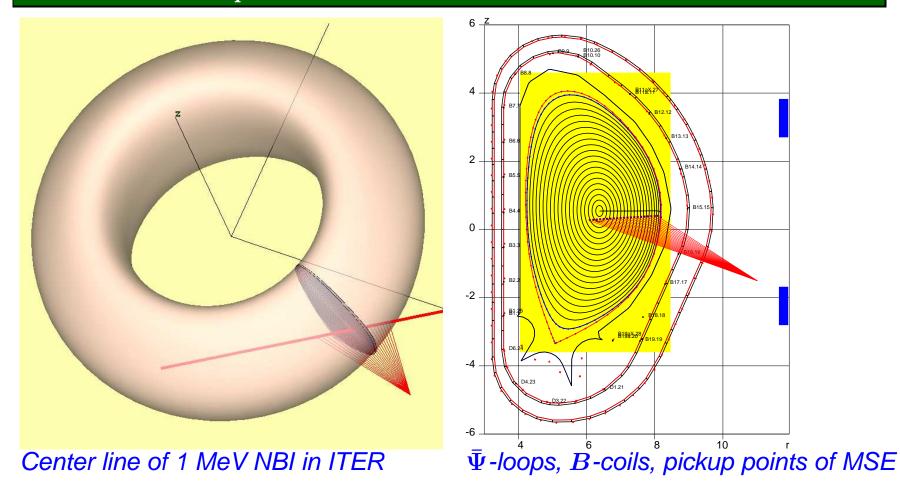


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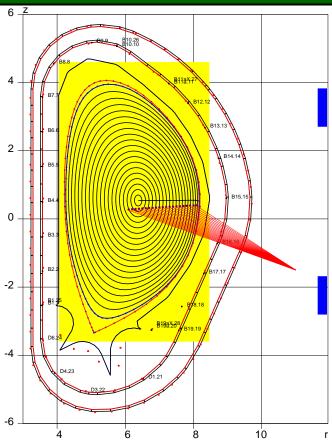
ITER B=5.3 T, I_{pl} =15 MA eta=2.8% equilibrium configuration



One of unique features of ITER is its 1 MeV neutral beam injection



Measurements of the Line Shift due to MSE was proposed by Nova Photonics as a diagnostics of ITER configuration



Reference signal errors ϵ used here for calculating variances in equilibrium reconstruction in ITER:

Signal name	$\epsilon^{relative}$	$\epsilon^{absolute}$	Comment
B-coils	0.01	0.01 T	local probes
Ψ -loops	0.01	0.001 Vsec	
Φ -loop	0.01	0.001 Vsec	diamagnetic loop
MSE-LP	0.01	0.1 ^o	B_z/B_{arphi} from MSE line polarization
MSE-LS	0.01	0.05 T	$\sqrt{ \mathbf{B} ^2 - (\mathbf{B} \cdot \mathbf{v})^2}$ from MSE line shift

MSE-LP and MSE-LS signals were assumed to be pointwise. This requires more realistic model from Nova Photonics.

The capabilities of equilibrium reconstruction with such a set of signal is the topic of the talk



The practice of equilibrium reconstruction (EqR) neglects making analysis of variances in reconstructed equilibiria

In tokamaks the Grad-Shafranov (GSh) equation describes the configuration

$$\Delta^*ar{\Psi} = -T(ar{\Psi}) - P(ar{\Psi})r^2, \quad T \equiv ar{F}rac{dar{F}}{dar{\Psi}}, \quad P \equiv \mu_0rac{dp}{dar{\Psi}}, \qquad (2.1)$$

Its solution can be perturbed: (a) by perturbation of the plasma shape

$$\xi(a_{pl},l), \tag{2.2}$$

and (b) by perturbation of two 1-D functions

$$\delta T(\bar{\Psi}), \quad \delta P(\bar{\Psi}).$$
 (2.3)

The question, neglected by present practice, is what level of perturbations cannot be distinguished given the finite accuracy of measurements.

The level of variances ξ , δT , δP determines the very value of reconstruction and of the entire diagnostics system



The theory of variances has been created in 2006 by L.Zakharov, J.Levandowski, V.Drozdov and D.McDonald

The problem is reduced to solving the linearized equilibrium problem

$$ar{\Psi} = ar{\Psi}_0 + \psi, \quad \Delta^* \psi + T'_{ar{\Psi}} \psi + P'_{ar{\Psi}} \psi = -\delta T(a) - \delta P(a) r^2$$
 (2.4)

for N possible perturbations

$$\xi = \sum\limits_{n=0}^{n < N_{\xi}} A_{n} \xi^{n}, \quad \delta T = \sum\limits_{n=0}^{n < N_{J}} T_{n} f^{n}, \quad \delta P = \sum\limits_{n=0}^{n < N_{P}} P_{n} f^{n},$$
 (2.5)

$$N = N_{\xi} + N_J + N_P, \quad f^{2n} = \cos 2\pi n a^2, \quad f^{2n+1} = \sin 2\pi n a^2,$$

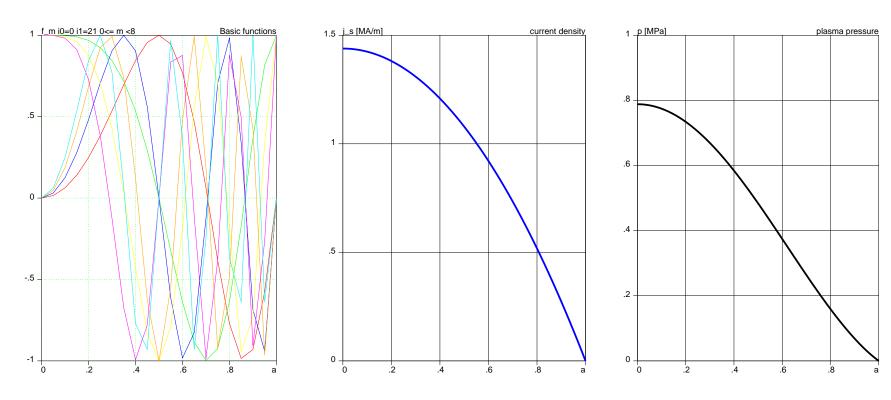
where $\xi^n(l)$, and $0 \le a \le 1$ is the square root from the normalized toroidal flux.

The response of the diagnostics to each of N solutions ψ^n can be calculated in a straightforward way.

ESC is based on linearization of the GSh equation. It was complemented with a routine for analysis of variances



8 functions $f^n(a^2)$ has been used to perturb $P(ar\Psi), T(ar\Psi)$



functions $f^n(a^2)$

sity profiles $\bar{j}_s(a)$

Trigonometric expansion background current den- background pressure profile $ar{p}(a)$

ESC can use an extended set of basis functions



After solving the perturbed GSh equation, the problem is reduced to a matrix problem

Let vector \vec{X} contains the amplitudes of perturbations

$$\vec{X} \equiv \left\{ \underbrace{A_0, A_1, \dots, A_{N_b-1}}_{N_\xi \text{ of } \xi}, \underbrace{T_0, \dots, T_{N_T-1}}_{N_T \text{ of } \delta T}, \underbrace{P_0, \dots, P_{N_P-1}}_{N_P \text{ of } \delta P}, \right\}$$
(3.1)

and vector $\delta \vec{S}$ represents the signals

$$\delta ec{S} \equiv \left\{ egin{aligned} \delta \Psi_0, \dots, \delta \Psi_{M_\Psi-1}, \delta B_0, \dots, \delta B_{M_B-1}, \delta S_0, \dots, \delta S_{M_S-1} \ M_\Psi \ \emph{of} \ \delta \Psi \end{aligned}
ight., \left\{ egin{aligned} \delta W_{0}, \dots, \delta W_{M_B-1}, \delta S_0, \dots, \delta S_{M_S-1} \ M_S \ \emph{of} \ \delta O \end{array}
ight. \end{cases}
ight.$$

32 Ψ , 1 $\Phi_{diamagnetic}$ -loops, 64 B-probes, 21 MSE-LP (line polarization) and 21 MSE-LS (line shift) signals (both pointwise) were used in the analysis.

ESC calculates the response matrix A relating $\delta \vec{S}$ and perturbations $\delta \vec{X}$

$$\delta \vec{S} = \mathsf{A} \vec{X}, \quad \mathsf{A} = \mathsf{A}_{M imes N}.$$
 (3.3)



The working matrix $\overline{\mathbb{A}}$ weights δS_m based on their accuracy

$$(\overline{\mathsf{A}})_m^n = \frac{1}{\epsilon_m} (\mathsf{A})_m^n, \quad \delta \bar{S}_m = \frac{1}{\epsilon_m} \delta S_m, \quad \overline{\mathsf{A}} \vec{X} = \delta \vec{\bar{S}},$$
 (3.4)

where ϵ_m is the error in the signal S_m . SVD expresses the matrix \overline{A} as a product

$$\overline{\mathbf{A}} = \mathbf{U} \cdot \mathbf{W} \cdot \mathbf{V}^{T},
\mathbf{U} = \mathbf{U}_{M \times N}, \quad \mathbf{U}^{T} \cdot \mathbf{U} = \mathbf{I}, \quad \mathbf{I}_{m}^{n} = \boldsymbol{\delta}_{m}^{n},
\mathbf{W} = \mathbf{W}_{N \times N}, \quad \mathbf{W}_{k}^{n} = \mathbf{w}^{n} \boldsymbol{\delta}_{k}^{n},
\mathbf{V} = \mathbf{V}_{N \times N}, \quad \mathbf{V}^{T} \cdot \mathbf{V} = \mathbf{I}.$$
(3.5)

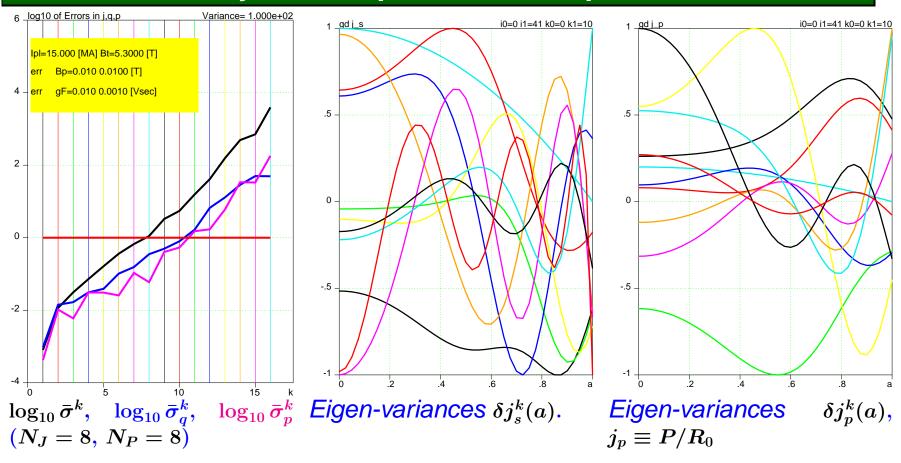
Here, w^n are the eigenvalues of the matrix problem.

The resulting vector of variances can be represented as a linear combination of "eigenvectors", which are the columns of matrix V

$$ec{X}^k = ec{V}^k, \quad \mathsf{A}ec{X}^k = w^k ec{U}^k, \quad ar{\sigma}^k \equiv \sqrt{rac{1}{M} \sum\limits_{m=0}^{m < M} \left(\mathsf{A}ec{X}^k
ight)_m^2} = rac{w^k}{\sqrt{M}}, \qquad (3.6)$$

Eq.(3.6) gives variances and normalized RMS $\bar{\sigma}^k$ in an explicit form. The perturbations \vec{X}^k with $\bar{\sigma}^k>1$ are "invisible" for diagnostics

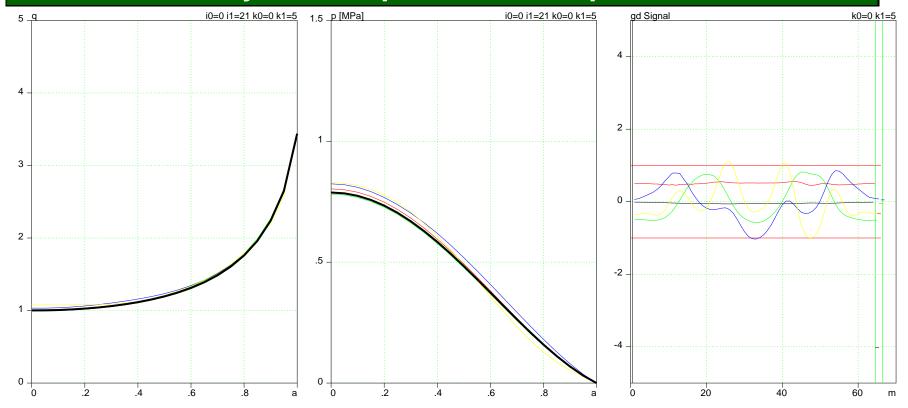




 $\bar{\sigma}_q$ and $\bar{\sigma}_p^k$ [MPa] on the left plot are RMS for q- and p-profiles

Perturbations $j_s^{k>8}, j_p^{k>8}$ are invisible and cannot be reconstructed



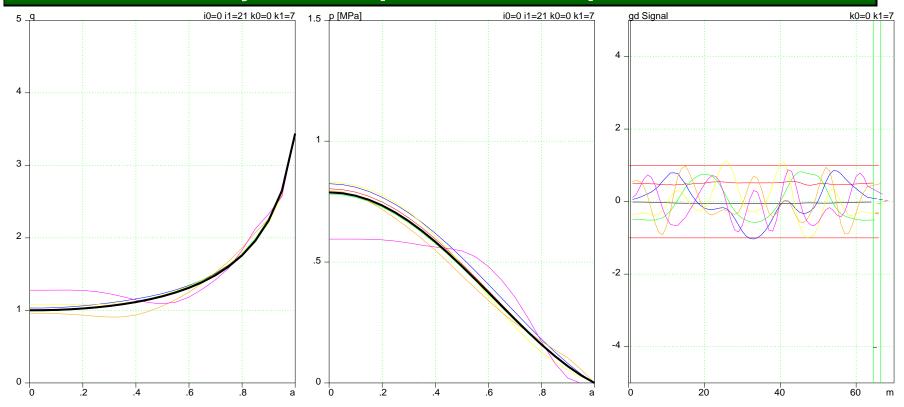


for $k_J \leq 3$, $k_P \leq 2$

q- profile and variances variances in p-profile as Signals $\delta S_m/\epsilon_m$ generated by perturbations functions of a

For $k_J+k_P=5$, typically used, the reconstruction looks very good KiloGb's of reconstructions "data" can be easily generated

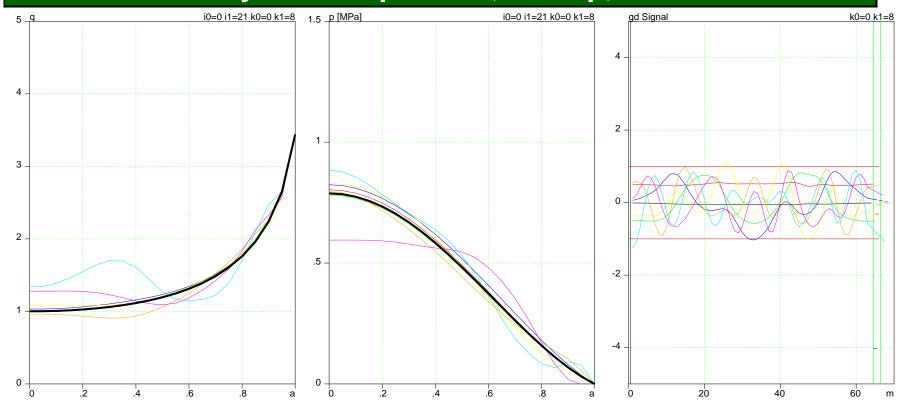




q- profile and variances p- profile and its vari- Signals $\delta S_m/\epsilon_m$ generfor $k_J \le 4$, $k_P \le 3$. ances as functions of a ated by perturbations

Testing $k_J + k_P = 7$ shows that the reconstruction is, in fact, not so good



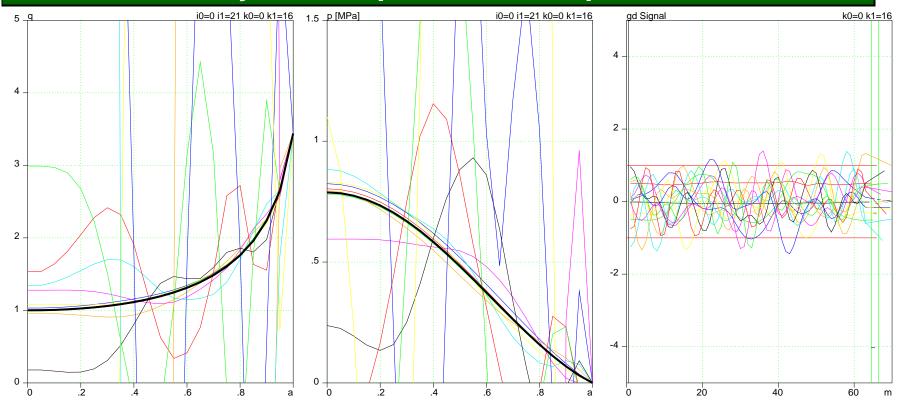


for $k_J < 4$, $k_P < 4$

q- profile and variances p- profile and its vari- Signals $\delta S_m/\epsilon_m$ generances as functions of a ated by perturbations

Testing $k_J + k_P = 8$ shows that even the q reconstruction is doubtful





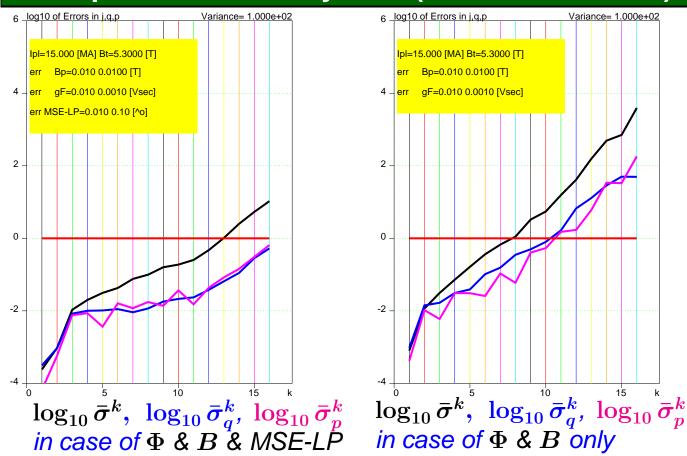
for $k_J \leq 8$, $k_P \leq 8$

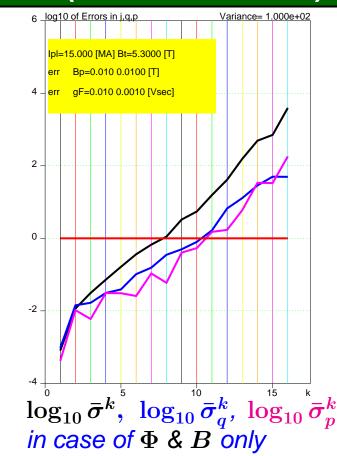
q- profile and variances p- profile and its vari- Signals $\delta S_m/\epsilon_m$ generances as functions of a ated by perturbations

Test of $k_J+k_P=16$ shows that with no constrains the reconstruction has no scientific value and is a sort of "beliefs"



Fixed plasma boundary with (Φ & B & MSE-LP) signals

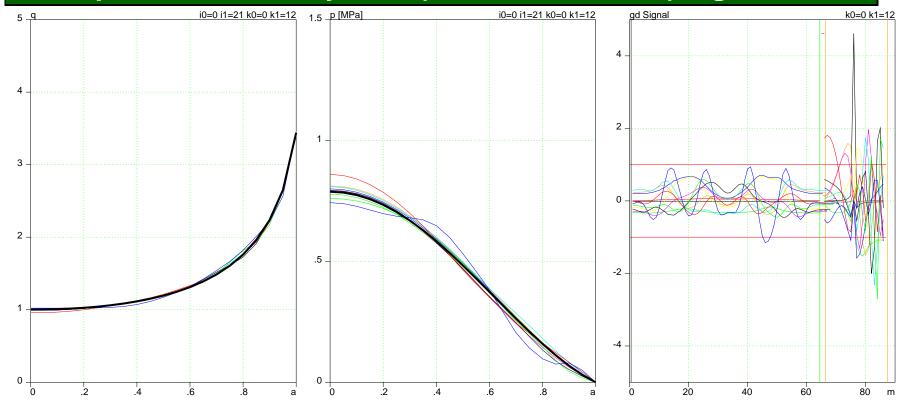




Use of MSE-LP drops largest RMS $\bar{\sigma}$, makes 12 perturbations visible, and dramatically improves reconstruction of q, p



Fixed plasma boundary with (Φ & B & MSE-LP) signals

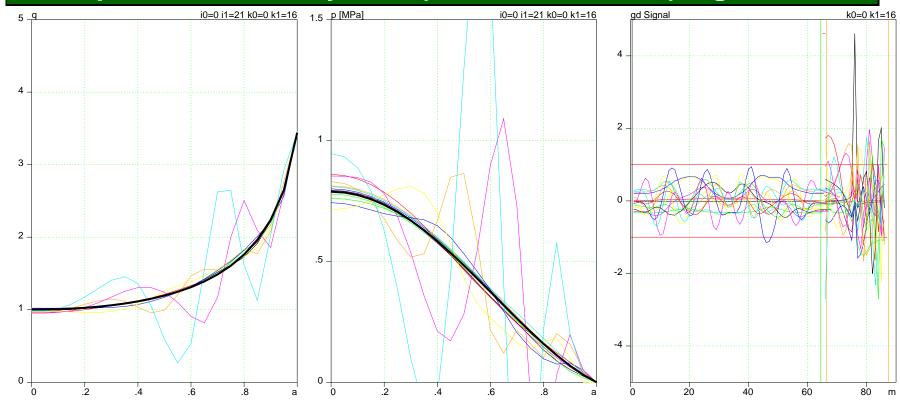


q- profile and variances p- profile and its vari- Signals $\delta S_m/\epsilon_m$ generated for $k_J \le 6$, $k_P \le 6$ ances as functions of a by perturbations

Testing N=12 shows that MSE-LP allows to reconstruct both q- and p-profiles



Fixed plasma boundary with (Φ & B & MSE-LP) signals



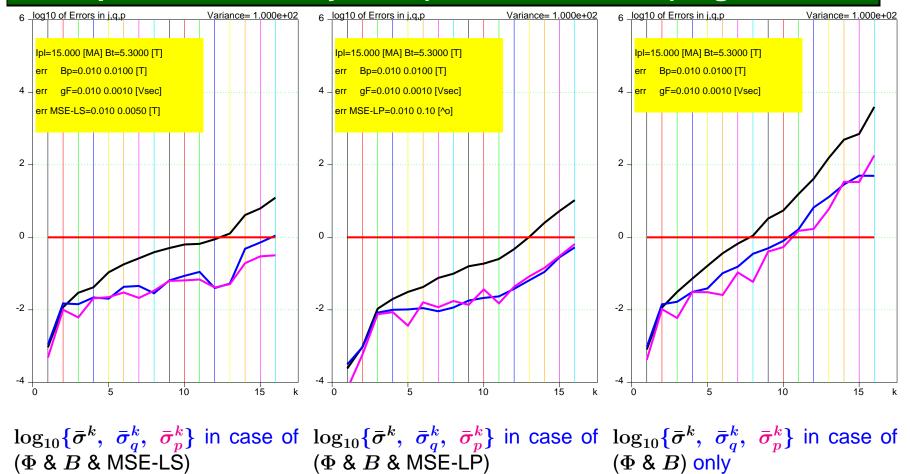
for all k

q- profile and variances p- profile and its vari- Signals $\delta S_m/\epsilon_m$ generances as functions of a ated by perturbations

Only perturbations with $k \geq 14$ might be potentially troublesome



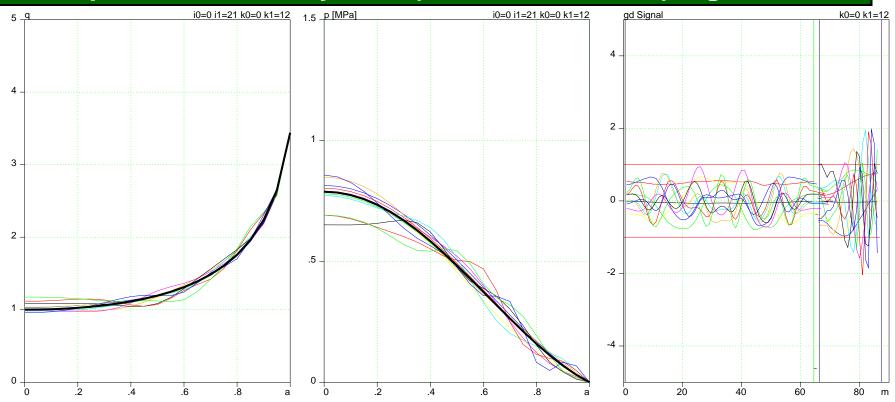
Fixed plasma boundary with (Φ & B & MSE-LS) signals



Use of MSE-LS can compete with MSE-LP in its value for reconstruction



Fixed plasma boundary with (Φ & B & MSE-LS) signals



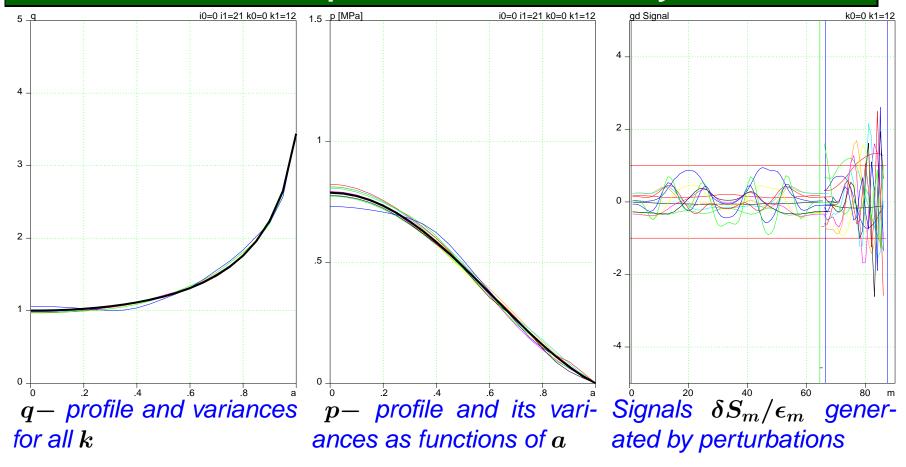
for all k

q- profile and variances p- profile and its vari- Signals $\delta S_m/\epsilon_m$ generances as functions of a ated by perturbations

Perturbations with $k \le$ 12 can be reconstructed using MSE-LS



Same case with the improved relative accuracy of MSE-LS

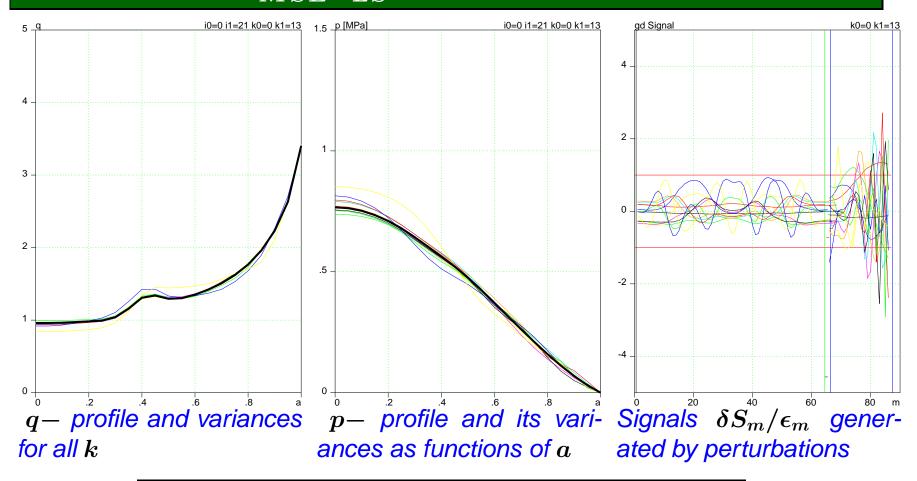


A realistic reduction of relative error $\epsilon_{MSE-LS}^{relative}
ightarrow 0.1\%$ improves

the pressure profile reconstruction



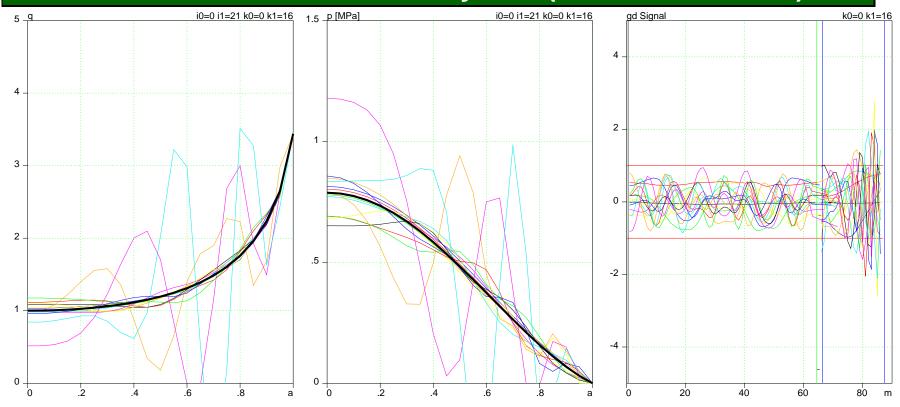
Same case with $\epsilon_{MSE-LS}^{relative} ightarrow 0.1\%$ and non-monotonic \overline{j}_s



MSE-LS can pick up the details of the current drive



Back to reference fixed boundary and (Φ & B & MSE-LS)

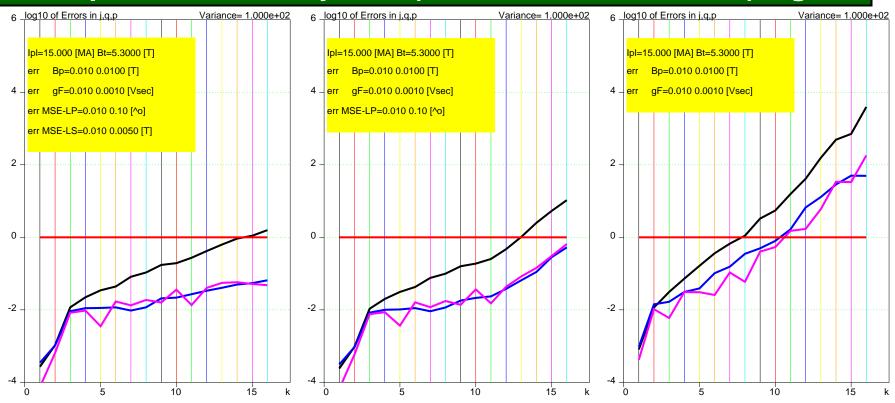


q- profile and variances p- profile and its vari- Signals $\delta S_m/\epsilon_m$ gener-for all k ances as functions of a ated by perturbations

With MSE-LS only perturbations with $k \ge 13$ might be potentially troublesome



Fixed plasma boundary with (Φ & B & MSE-LP & -LS) signals

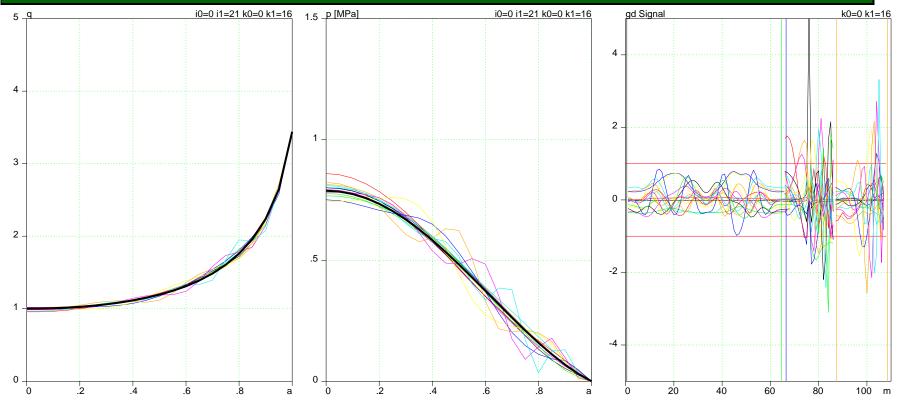


 $\log_{10}\{\bar{\sigma}^k,\ \bar{\sigma}^k_q,\ \bar{\sigma}^k_p\} \text{ in case of } \log_{10}\{\bar{\sigma}^k,\ \bar{\sigma}^k_q,\ \bar{\sigma}^k_p\} \text{ in case of } \log_{10}\{\bar{\sigma}^k,\ \bar{\sigma}^k_q,\ \bar{\sigma}^k_p\} \text{ in case of } (\Phi \& B \& \text{MSE-LP}) \qquad \qquad (\Phi \& B) \text{ only}$

Both MSE-LP & LS allows for a reliable reconstruction of q- and p-profiles



Fixed plasma boundary with (Φ & B & MSE-LP&LS) signals



for all k

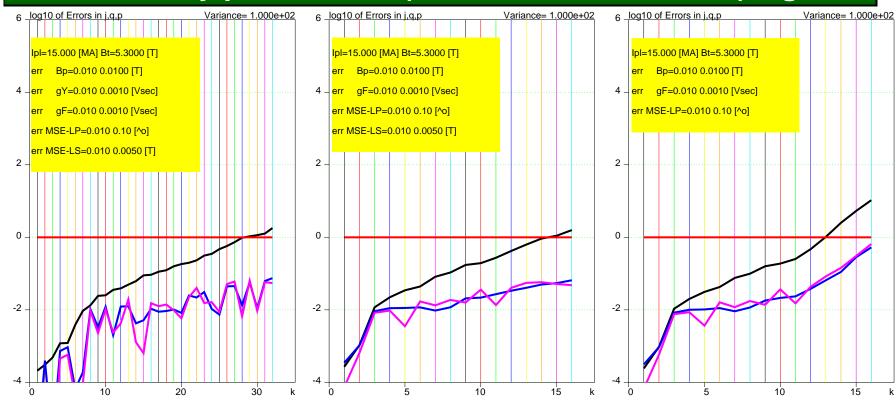
as functions of a

q- profile and variances p-profile and its variances Signals $\delta S_m/\epsilon_m$ generated by perturbations

q- and p-profiles can be reconstructed in all spectrum of k



Free boundary plasma with (Φ & B & MSE-LP & -LS) signals

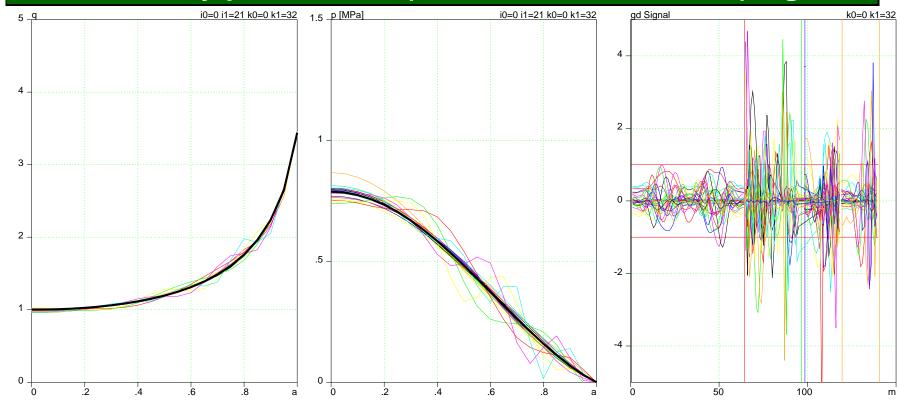


$$\begin{array}{ll} \log_{10}\{\bar{\sigma}^k,\ \bar{\sigma}^k_q,\ \bar{\sigma}^k_p\} \ \text{in case of} \ \log_{10}\{\bar{\sigma}^k,\ \bar{\sigma}^k_q,\ \bar{\sigma}^k_p\} \ \text{in case of} \\ (\Phi \& B \& \text{MSE-LP} \& \text{MSE-LS}), \ \ (\Phi \& B \& \text{MSE-LP} \& \text{MSE-LS}), \\ \vec{\xi} \neq 0 & \vec{\xi} = 0 \end{array} \\ \begin{array}{ll} \log_{10}\{\bar{\sigma}^k,\ \bar{\sigma}^k_q,\ \bar{\sigma}^k_p\} \ \text{in case of} \\ (\Phi \& B \& \text{MSE-LP}), \ \vec{\xi} = 0 \end{array}$$

Free boundary expands the k range but does not affect the reconstruction



Free boundary plasma with (Φ & B & MSE-LP & -LS) signals



for all extended k

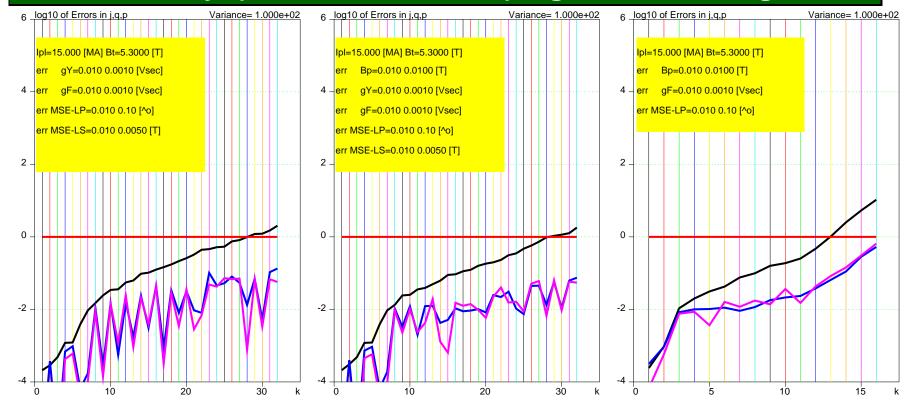
as functions of a

q- profile and variances p-profile and its variances Signals $\delta S_m/\epsilon_m$ generated by perturbations

q- and p-profiles can be reconstructed in all extended spectrum of k



Free boundary, (Φ & MSE-LP & -LS) signals, NO B-signals

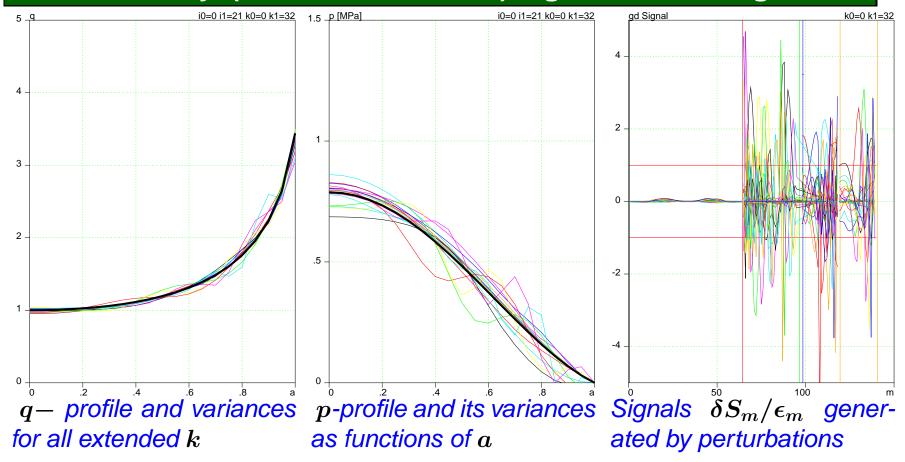


 $\begin{array}{ll} \log_{10}\{\bar{\sigma}^k,\ \bar{\sigma}^k_q,\ \bar{\sigma}^k_p\} \ \text{in case of} \ \log_{10}\{\bar{\sigma}^k,\ \bar{\sigma}^k_q,\ \bar{\sigma}^k_p\} \ \text{in case of} \\ (\Phi \& B \& \text{MSE-LP \& MSE-LS}), \ (\Phi \& B \& \text{MSE-LP \& MSE-LS}), \\ \bar{\xi} \neq 0 \ \qquad \qquad \bar{\xi} \neq 0 \end{array} \\ \begin{array}{ll} \log_{10}\{\bar{\sigma}^k,\ \bar{\sigma}^k_q,\ \bar{\sigma}^k_p\} \ \text{in case of} \\ (\Phi \& B \& \text{MSE-LP}), \ \text{and} \ \bar{\xi} = 0 \end{array}$

(MSE-LP & MSE-LS) together can do the job for external B-coils



Free boundary, (Φ & MSE-LP & -LS) signals, NO B-signals



q- and p-profiles can be reconstructed over extended spectrum of k even with NO B-coil signals



The capability of calculating variances, now developed, has completed the theory of equilibrium reconstruction

- 1. The quantitative evaluation of the quality of diagnostics systems can be done based on spectrum of "visible" perturbations
- 2. It was confirmed that the internal measurements of the magnetic field are crucial for reconstruction
- 3. Either MSE-LP (line polarization) or MSE-LS (line shift) signals from the plasma in addition to external measurements allow for a complete reconstruction (of both q- and p-profiles).
- 4. The presented technique can be used to optimize the MSE diagnostic set plans on any tokamaks.
- 5. The proposal by Nova Photonics to utilize the high magnetic field and NBI energy in ITER for extraction of MSE-LS signals would significantly enhance the equilibrium reconstruction capability in ITER.

The extension of the theory should be focused on realistic simulation of signals used in reconstructions

